

# Scattering diagram of $(\mathbb{P}^2, E)$ part II

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Based on work of

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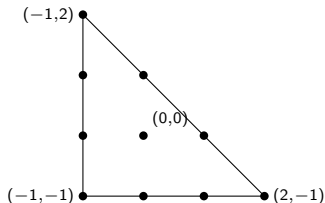
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## A log smooth degeneration of $(\mathbb{P}^2, E)$

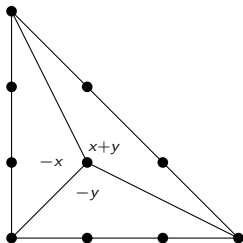
### First steps

- Degenerate  $\mathbb{P}^2 = \mathbb{P}_\Delta \rightsquigarrow \cup_i \mathbb{P}_{\Delta_i}$
- $\mathbb{P}_\Delta \supset E \rightsquigarrow D = \cup_i D_i \subset \cup_i \mathbb{P}_{\Delta_i}$ ,  $D_i \subset \mathbb{P}_{\Delta_i}$
- $\cup_i \Delta_i$  is a regular polyhedral decomposition of  $\Delta$

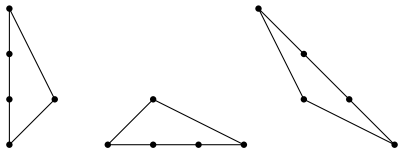
Let  $M = \mathbb{Z}^2$  and consider the polytope  $\Delta$  for  $(\mathbb{P}^2, \mathcal{O}(3))$ :



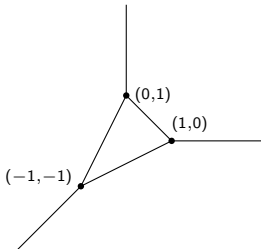
Take an integral regular subdivision of  $\Delta$ , i.e. a subdivision induced as the “bending locus” of a piecewise linear convex function  $\varphi$ :



This determines a degeneration of  $\mathbb{P}^2$  into the union of three weighted projective spaces  $\mathbb{P}(1, 1, 3) = \mathbb{P}_{\Delta_i}$ , where  $\Delta_i$  is either of

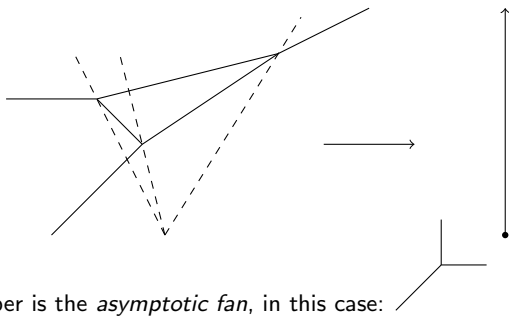


To describe the degeneration of  $\mathbb{P}^2$  at the fan level, start with



The “local fans” at the vertices are the fans for the  $\mathbb{P}(1, 1, 3)$ ’s.

Move the polyhedral decomposition of  $\mathbb{R}^2$  to “height” 1 and take the cone over it, as well as the natural map to  $\mathbb{A}^1$  given by “height”:



The generic fiber is the *asymptotic fan*, in this case:

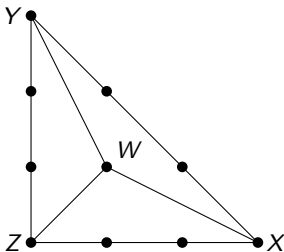
The components of the central fiber are given by the “quotient” by the vertical rays  $\sim \mathbb{P}_{\Delta_i} \cong \mathbb{P}(1, 1, 3)$ .

This is the desired degeneration  $\mathbb{P}_{\Delta} \rightsquigarrow \cup_i \mathbb{P}_{\Delta_i}$ . (More projectively, the total family can also be described from the dual non-compact polyhedron.)

In this degeneration the toric divisor of  $\mathbb{P}^2$  goes to a cycle of lines, one in each of  $\mathbb{P}_{\Delta_i}$ .

## Equations

Integral points of the polytope correspond to sections of ample line bundles of the family:



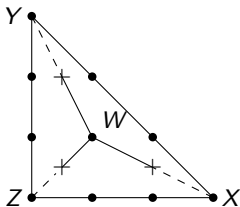
$\varphi$  is defined by taking value 1 at  $X, Y, Z$  and value 0 at  $W$ . The smoothing parameter is  $t = z^{(0,0,1)}$ . This leads to the family

$$V(XYZ - t^3 W) \subset \mathbb{P}(1, 1, 1, 3) \times \mathbb{A}^2$$

$V(W)$  is the toric horizontal divisor.

## Gross-Siebert generalization

Integral points of the Gross-Siebert base correspond to sections of ample line bundles of the family:



We add a parameter  $s$  that pushes in the corners.  $g$  is a general cubic polynomial (e.g.  $g = X^3 + Y^3 + Z^3$ ).

$$V(XYZ - t^3(W + sg)) \subset \mathbb{P}(1, 1, 1, 3) \times \mathbb{A}^2$$

$V(W)$  is the horizontal divisor. It is a degeneration of a smooth cubic into a union of lines.

The family has 3 log singularities on smooth points of the common divisor of each of two  $\mathbb{P}(1, 1, 3)$ . Blowing them up we get a log smooth family.

## Outline of proof of Gräfnitz

- Consider the just constructed log smooth degeneration of  $(\mathbb{P}^2, D)$ .
- Apply the decomposition formula of Abramovich-Chen-Gross-Siebert, indexed by tropical curves  $h$  in the dual intersection complex.
- Each term is a weighted count of curves that tropicalize to  $h$ .
- Compute the latter by applying GPS locally.
- Reformulate inside the scattering diagram, which is the dual intersection complex of the degeneration.

### Conclusion

The wall-crossing functions encode  $N_{0,d}^k$ .

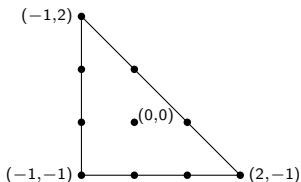


## Relative/log Gromov-Witten invariants

- $P_d := 3d$ -torsion points of  $E$ .
- For  $p \in P_d$ ,  $d(p) := \min\{k \mid p \in P_k\}$ .
- $\overline{M}_0(\mathbb{P}^2/E, d) :=$  moduli space of degree  $d$  genus 0 basic stable log maps of maximal tangency (J. Li, Chen, Abramovich-Chen, Gross-Siebert)
- $\overline{M}_0(\mathbb{P}^2/E, d) = \bigsqcup_{p \in P_d} \overline{M}_0(\mathbb{P}^2/E, d)^p$
- $N_{0,d} := \int_{[\overline{M}_0(\mathbb{P}^2/E, d)]^{\text{vir}}} \mathbf{1} \in \mathbb{Q}$
- $N_{0,d}^p := \int_{[\overline{M}_0(\mathbb{P}^2/E, d)^p]^{\text{vir}}} \mathbf{1} \in \mathbb{Q}$
- $N_{0,d} = \sum_{p \in P_d} N_{0,d}^p$
- $N_{0,d}^p$  depends only on  $d(p)$
- For  $k|d$ ,  $N_{0,d}^k := N_{0,d}^p$  for  $p \in P_d$  with  $d(p) = k$

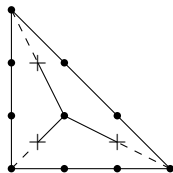
# Gross-Siebert construction of dual intersection complex

## Polytope of $(\mathbb{P}^2, \mathcal{O}(3))$

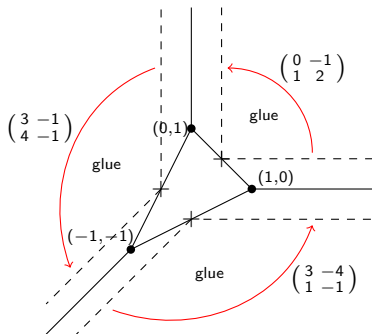


## Smooth out boundary

by pushing singularities inside,  
creating 3 focus-focus singularities



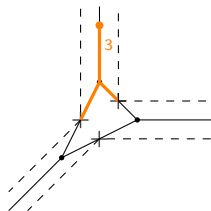
## Dual intersection complex



The monodromy around each singularity is conjugate to  $\begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$

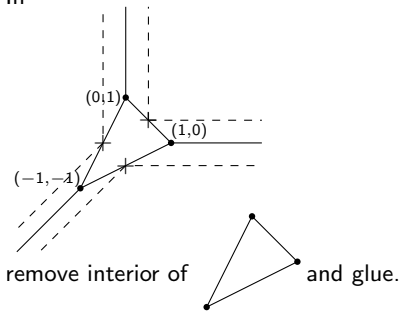
Each singularity has a monodromy-invariant line along which disks propagate.

Disks satisfy the balancing condition:



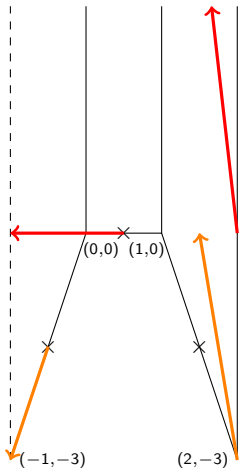
## Fundamental domain

In



Respect balancing condition at  $\bullet$ ,  
 only one infinite direction,  
 global monodromy is

$$\begin{pmatrix} 0 & -1 \\ 1 & 2 \end{pmatrix} \cdot \begin{pmatrix} 3 & -4 \\ 1 & -1 \end{pmatrix} \cdot \begin{pmatrix} 3 & -1 \\ 4 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -9 & 1 \end{pmatrix}$$



## Walls carry wall-crossing functions

Away from the discriminant locus,  $B$  admits an integral affine structure  $\Rightarrow$  sheaf of integral tangent vectors  $\Lambda_B$ .

$$\mathbb{C}[\Lambda_B \oplus \mathbb{Z}] := \langle z^{(m,a)} \mid m \in B, a \in \mathbb{Z} \rangle$$

$z^{(0,1)}$  is the smoothing parameter of the mirror family (to get the smoothing we need the multivalued piece-wise affine function  $\varphi$ , which we will continue to ignore).

### The original wall-crossing function

Let  $p$  be the wall directed by  $m_p \in \Lambda_B$  coming out of a focus-focus singularity. Its wall-crossing function is

$$\begin{array}{c} \times \\ \hline 1 + z^{(-m_p, 0)} \\ \hline \xrightarrow{m_p} \end{array}$$

### Proposition (Gräfnitz '20)

For a wall  $(p, m_p)$ , denote by  $\mathfrak{H}_{p,w}$  the set of *tropical disks* ending on  $p$  with final edge directed by  $-wm_p$ .

Each  $h \in \mathfrak{H}_{p,w}$  carries a combinatorial multiplicity  $N_h$ . Then

$$\log f_p = \sum_{w=1}^{\infty} \sum_{h \in \mathfrak{H}_{p,w}} N_h z^{(-wm_p, 0)}$$

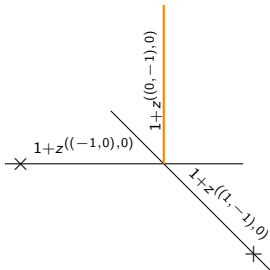
### Example

$$\frac{1 + z^{(-m_p, 0)}}{\times}$$

$$\log \left( 1 + z^{(-m_p, 0)} \right) = \sum_{w=1}^{\infty} \frac{(-1)^{w-1}}{w} z^{(-wm_p, 0)}$$

## Simple scattering

If two initial walls  $p, p'$  meet with  $\det(m_p, m_{p'}) = \pm 1$ , then only one additional disk is produced directed by  $m_p + m_{p'}$  and with wall-crossing function  $1 + z^{(-m_p - m_{p'}, 0)}$ :



If  $\det(m_p, m_{p'}) = \pm 2$ , there is infinite scattering with rays directed by  $n m_p + (n + 1)m_{p'}$ ,  $(n + 1)m_p + n m_{p'}$  and  $m_p + m_{p'}$ .

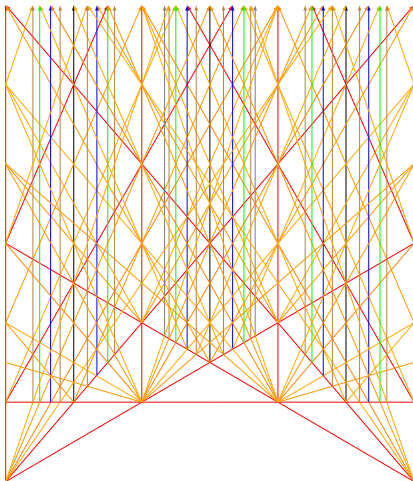
## General scattering

If  $\det(m_p, m_{p'}) = \pm 3$ , there are two regions with countably infinite scattering and a region where all rays with rational slope occur.

In general (also with more than 2 incoming rays), the scattering algorithm is described by GPS with coefficients given as log GW invs.

It works  $z^{(0,1)}$ -order by  $z^{(0,1)}$ -order.

Consistency condition needs to be satisfied at each joint.





- The circle at  $\infty$  is  $\mathbb{R}/3\mathbb{Z}$  with the vertical rays ending in  $\mathbb{Q}/3\mathbb{Z}$ .
- $G$  abelian group,  $x \in G$  with  $3 \mid \text{ord}(x)$ ,  $d(x) := \min\{d \mid 3d x = 0\}$
- $r_\ell := \#\{x \in \mathbb{Z}/3\ell \mid d(x) = \ell\}$
- $s_{k,\ell} := \#\{(a, b) \in \mathbb{Z}/3k \times \mathbb{Z}/3k \mid d((a, b)) = k, d(a) = \ell\}$

$x \in \mathbb{Q}/3\mathbb{Z}$ ,  $d(x) = \ell$ , then

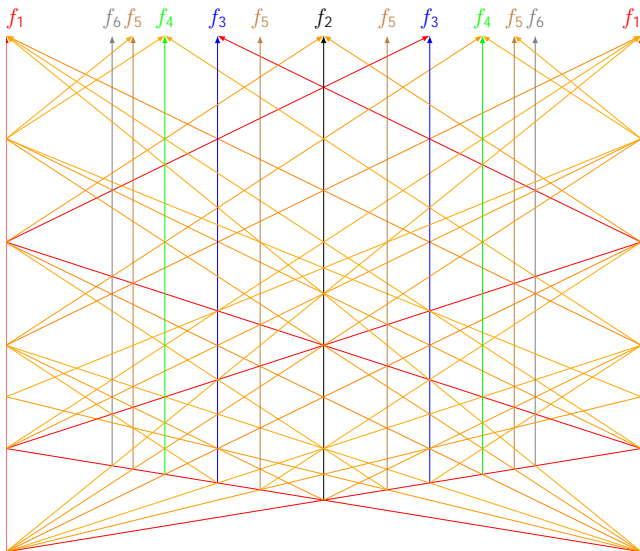
$$\frac{s_{k,\ell}}{r_\ell} = \#\{p \in P_d \mid d(p) = k, p \text{ specializes to } x\}$$

### Theorem (Gräfnitz '20)

For  $x = z^{((0,-1),0)}$  and  $p$  an unbounded wall of order  $\ell$

$$\log f_p = \sum_{d=1}^{\infty} 3d \left( \sum_{\ell \mid k \mid d} \frac{s_{k,\ell}}{r_\ell} N_{0,d}^k \right) x^{3d}.$$

## Zoom into the central third (image from Gräfnitz)



## Also from Gräfnitz

The sage code gives the following:

$$\log f_1 = 9x^3 + \frac{63}{2}x^6 + 246x^9 + \frac{9279}{4}x^{12} + \frac{175464}{5}x^{15} + 307041x^{18} + O(x^{21})$$

$$\log f_2 = 36x^6 + 2322x^{12} + 307164x^{18} + O(x^{21})$$

$$\log f_3 = 243x^9 + \frac{614061}{2}x^{18} + O(x^{21})$$

$$\log f_4 = 2304x^{12} + O(x^{21})$$

$$\log f_5 = 25425x^{15} + O(x^{21})$$

$$\log f_6 = 307152x^{18} + O(x^{21})$$

## Insight of Bousseau

Using

$$\log f_p = \sum_{d=1}^{\infty} 3d \left( \sum_{\ell|k|d} \frac{s_{k,\ell}}{r_\ell} N_{0,d}^k \right) x^{3d},$$

how do you prove that the BPS invariant  $\Omega_{d,k}^{p^2/E}$  is independent of  $k$ ? Here

$$(-1)^{d-1} N_{0,d}^k = \sum_{k|d'|d} \frac{1}{(d/d')^2} \Omega_{d',k}^{p^2/E}.$$

### Idea

Translate the question to a question about DT invariants of  $K_{p^2}$ .

# Heuristics I

# Heuristics II

# Heuristics III

$$\begin{aligned}
 U &:= \left\{ (x, y) \in \mathbb{R}^2 \mid y > -\frac{x^2}{2} \right\} & U &\hookrightarrow \text{Stab } D^b(\mathbb{P}^2), \quad \sigma \mapsto (\mathcal{A}^\sigma, Z^\sigma) \\
 K_0(\mathbb{P}^2) &:= \Gamma \simeq \mathbb{Z}^3 \ni (r(F), d(F), \chi(F)) & \text{Central charge } Z^\sigma &: \Gamma \rightarrow \mathbb{C}, \quad \gamma \mapsto Z_\gamma^\sigma \\
 & & F \in \mathcal{A}^\sigma, \text{ phase } \phi(F) &:= \frac{1}{\pi} \arg Z_{\gamma(F)}^\sigma
 \end{aligned}$$

$$\sigma \in U, \gamma \in \Gamma \rightsquigarrow M_\gamma^\sigma := \{F \mid \sigma\text{-semistable}, \gamma(F) = \gamma\} / S\text{-equivalence}$$

Wall and chamber structure on  $U$ ,  $M_\gamma^\sigma$  jumps along walls. For  $y \gg 0$ ,  $(x, y)$ -semistability  $\rightsquigarrow$  Gieseker-semistability.

$$\Omega_\gamma^\sigma := (-1)^{\dim M_\gamma^\sigma} \text{le}(M_\gamma^\sigma) \in \mathbb{Z}$$

satisfy Kontsevich-Soibelman wall-crossing.



## Wall-structure

For  $\gamma \in \Gamma$ , rays

$$R_\gamma := \{\sigma \in U \mid \operatorname{Re} Z_\gamma^\sigma = 0, \Omega_\gamma^\sigma \neq 0\}$$

$\leadsto$  straight line in direction  $(-r, d)$ .  $R_\gamma$  decomposes into segments, where the  $\Omega_{k\gamma}^\sigma$  do not jump. To these we attach the function

$$\sum_{k \geq 1} \frac{\Omega_{k\gamma}^\sigma}{k^2} z^{(kr, -kd)}.$$

### Theorem (Bousseau '19)

This defines a consistent scattering diagram.

## Comparing the two scattering diagrams

### Theorem (Bousseau '19)

Both scattering diagrams have the same initial conditions  $\Rightarrow$  they are the same.

## Going to $\infty$

### Theorem (Bousseau '19)

The function attached to unbounded vertical rays with  $\gamma = (0, d, \chi)$  is

$$(-1)^{d-1} \left( \sum_{\ell|(d,\chi)} \frac{1}{\ell^2} \Omega_{\frac{d}{\ell}, \frac{\chi}{\ell}}^{p^2} \right) z^{(0,-d)}$$

### Corollary

Let  $\ell_{d,\chi} := \frac{d}{(d,\chi)}$ . Then

$$(-1)^{d-1} \sum_{\ell|(d,\chi)} \frac{1}{\ell^2} \Omega_{\frac{d}{\ell}, \frac{\chi}{\ell}}^{p^2} = \sum_{\ell_{d,\chi} | k | d} \frac{s_{k,\ell_{d,\chi}}}{r_{\ell_{d,\chi}}} N_{0,d}^k$$

Moreover,

$$\sum_{\ell|(d,\chi)} \frac{1}{\ell^2} \Omega_{\frac{d}{\ell}, \frac{\chi}{\ell}}^{p^2} = \sum_{\ell_{d,\chi} | d' | d} \frac{1}{(d/d')^2} \Omega_{d', \chi \frac{d'}{d}}^{p^2}$$

Hence

$$\sum_{\ell_{d,\chi} | d' | d} \frac{1}{(d/d')^2} \Omega_{d',\chi}^{p^2} = \sum_{\ell_{d,\chi} | k | d} \frac{s_{k,\ell_{d,\chi}}}{r_{\ell_{d,\chi}}} (-1)^{d-1} N_{0,d}^k$$

**Corollary (Bousseau '19)**

$$\Omega_{d,\chi}^{p^2} = \sum_{\ell_{d,\chi} | k | d} \frac{s_{k,\ell_{d,\chi}}}{r_{\ell_{d,\chi}}} \Omega_{d,k}^{p^2/E}$$

Conclude by using

**Theorem (Bousseau '19)**

$\Omega_{d,\chi}^{p^2}$  is independent of  $\chi$ .

Thank you